NEW FIR LASER LINES AND FREQUENCY MEASUREMENTS FOR OPTICALLY PUMPED CD,OH

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Twenty new cw FIR laser lines in CD_3OH , optically pumped by a CO_2 laser, are reported. The frequencies of 39 of the stronger laser lines were measured relative to stabilized CO_2 lasers with a fractional uncertainty, as determined by the reproducibility of the FIR frequency itself, of 2 parts in 10^7 .

Key Words: CD₃OH laser lines, laser frequency measurements, optically pumped FIR laser, new laser lines, relative intensity, relative polarization.

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Introduction

The set of isotopes of methyl alcohol - CH $_3$ OH, CD $_2$ OH, CH $_3$ OD, CH $_2$ DOH, CHD $_2$ OH, and 13 CH $_3$ OH - occupies a prominent position in the list of cw optically pumped far infrared (FIR) lasing molecules, accounting for approximately 35% of all the presently tabulated lasing lines (1). Even more important is the fact that these isotopes are responsible for about 130 of the approximately 160 lines found at wavelengths shorter than 200 um. Hence, this series of isotopes alone constitutes the majority of all of the frequency coverage needed for frequency metrology and laser spectroscopy in this region of the FIR. This laboratory has been carrying out a program to systematically measure frequencies of known methanol lasing transitions and to search for new FIR laser lines. The results of frequency measurements, accurate to a few parts in 107, have been published for CH₃OH, CH₃DOH, and CD₃OD (2 - 4). In this paper we report similar measurements on 39 of the stronger lasing transitions in CD₃OH (methanol-d₃).

The first account of optically-pumped FIR lasing action in CD3OH was published by Dyubko et al. (5) who found 70 cw transitions ranging from 180 to 1300 µm. Subsequently, Danielewicz and Weiss (6) reported finding 21 new cw lines at shorter wavelengths, Grinda and Weiss (7) found another 10 new lines, Yoshida et al. (8) found 7 new lines, Sigg et al. (9) found a number of new short wavelength lines, and Pereira et al. (10) found 69 new lines, many of them at long wavelengths. In all cases, a CO2 laser was used to pump a conventional hole-coupled Fabry-Perot FIR resonator, and wavelengths were measured with an accuracy of about 0.5%. At 100 µm this corresponds to an uncertainty of 15 GHz in the laser frequency, which is entirely too crude for applications in frequency synthesis or spectroscopy. The frequency measurements reported in this paper correspond to over a thousandfold improvement in accuracy over the wavelength determinations.

Experimental

The CD₃OH lines were pumped by a 3 m, invarstabilized, grating- and piezoelectrically-tuned $\rm CO_2$ laser having single line cw powers of up to 30 w. The

CO2 laser power was coupled into a Fabry-Perot resonator $(R_1 = 200 \text{ cm}, R_2 = \infty, D = 100 \text{ cm})$ by focusing the beam through a 1 mm diameter hole in the center of the curved 12.5 cm diameter gold-coated copper mirror with 1 m focal length optics. FIR power was coupled out of the resonator with a 3 mm hole located 5.5 mm from the center of the same mirror. The position of the flat mirror at the other end of the cavity was adjusted with e micrometer drive to change the resonant frequency of the cavity. Irises in front of each mirrors were used to discriminate against higher order modes. Wavelengths of the laser lines accurate to ± 0.1 µm were measured by counting the number of modes over a 5 mm scan, using the known wavelength of the 118.8 µm line of CH3OH for calibration. This level of accuracy was needed to predict the lasing frequency to within the 1500 MHz detection bandwidth.

For the frequency measurements, the FIR laser output was focused onto a long-wire (6 mm), W-Ni, point-contact diode. The angle between the beam and the wire was adjusted to maximize the coupling. FIR frequencies were synthesized with two $^{12}\text{C}^{16}\text{O}_2$ lasers and an X-band klystron, as described previously (11). Each CO_2 laser was actively stabilized to the standing-wave saturation resonance in a low-pressure (5.3 Pa), intracavity, CO_2 absorption cell (12). The X-band klystron was stabilized to a quartz crystal, and its frequency was counted. The uncertainty in this synthesized FIR frequency was less than 100 kHz. The beat frequency between the synthesized frequency and the FIR laser frequency was amplified (25 db) and measured with a 1500-MHz spectrum analyzer and marker oscillator.

Results

In Table I, we reported 20 new laser lines from a CD₃OH sample of 99% isotopic purity. We have listed all of the lines reported in References (6) and (7) and those that oscillate in our laser from ref. (8), (9), and (10). Because of the same purity problems mentioned in Ref. (5), we have listed only those lines from that source which we could reproduce with our apparatus. Polarization of the FIR radiation relative to that of the CO₂ pump radiation was measured for about two-thirds of the listed lines.

Frequency measurements made for 39 lasers lines of CD₂OH are presented in Table II. The vacuum wavenumber is derived from the measured frequencies with c=299 792 458 m/s (13). Pressures given in the table for the various lines were measured with an air-calibrated thermocouple gauge and correspond to the maximum in the FIR laser output power. Polarizations of the lines relative to the linearly-polarized CO2 pump laser were measured with an external metal-mesh polarizer. Relative powers of the lines were determined by monitoring the laser output with a diamond window Golay cell with a 0.24-mm-thick crystal-quartz filter to block the CO2 radiation. Other calibrated attenuators were used with strong lasing lines to prevent saturation of the detector. Relative powers in Table I correspond to the maximum output of each line and are labelled only VS, S, M, W, or VW because the actual intracavity power varies with coupling, wavelength, and pump power, and no attempt was made to ascertain these dependencies.

Discussion of Results

The estimated fractional frequency uncertainty of about ± 2 x 10⁻⁷ results from a combination of factors. Principally, the accuracy is limited by our ability to set the free-running FIR laser to the center of its unperturbed gain-profile. This problem is compounded by shifts caused by competing lines and modes as well as by splittings and symmetries introduced into the lineshape by the ac-Stark and anisotropic-gain effects (14, 15). Variations up to several megahertz from these effects have occurred in past measurements but can usually be reduced by up to a factor of ten with great care in the measurement. Pressure shifts are small and are difficult to observe over the normal operating pressure range of the FIR laser.

Finally, we note that the 39 CD₃OH measured frequencies in this work are those for lines that oscillated easily in our system. Seventeen of these lines are reported here for the first time; the 22 others were observed in the earlier work. Of the 70 cw lines reported in Ref. (5), we find only 13 that oscillate in our apparatus; all lines reported in References (6) and (7) were observed. In addition, we were able to separate

the $R_{\rm I}(34)$ pump frequency into three separate offsets with the use of a FIR laser transversely pumped to eliminate frequency jumping at the ends of the tuning curve of ${\rm CO_2}$ caused by ${\rm CO_2}$ laser radiation feedback.

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TABLE I. CD3OH Lasing Lines Arranged by Pump

CO ₂ Pump Line	CD ₃ OH Laser Line λ(μm)	Relative Polarization	CD ₃ OH Power	Ref.
9.4 Band	VI			
R _{II} (34)	54.0 60.8	#	VS VS	7
R _{II} (28)	40.0 49.8 55.6 159.4 ^a 181.7 ^a 370.5 ^a	i H H	VS S M M	new 7 new 7 8 5
R _{II} (26)	50 498.7	<u>.</u>	-	new
R _{II} (22)	583	¥	W	5,7
R _{II} (14)	120.7 ^a 182.6 ^a 236 352.5 ^a	# # #	M M M M	8 7 7 7
R _{II} (6)	35.5 48.4 56.9	<u> </u>	W M W	new 9 7
P _{II} (8)	223	τ .	W	7
P _{II} (28)	435	1	VW	7
P _{II} (32)	351	K	VW	7

R_I (24) 61.5

Table I	(Continued)			
P _{II} (36)	189.7ª	11	S	10
P _{II} (40)	198.7 ^a 286.2 ^a	 4	M W	7 5

111(20)	103.7	H	3	10
P _{II} (40)	198.7 ^a 286.2 ^a	II R	M W	7 5
10.4 µm	Band			
R _I (45)	68.8	-	-	new
R _I (38)	50 122.2ª	Ħ	W M	new
R _I (36)	253.7ª 418.7ª	T d	S M	6 5,6
R) (34)'	43.1 168.1 ^a 180.7 ^a 222.2 _a 386.0 _a 430.9 ^a	_ T T	S S S - W	9 new 6 new 10
R ^b (34)"	86.7ª 264.8ª	Ţ Ţ	M M	6 5,6
Rb (34)"	'128.0 ^a 191.4 ^a	1	S M	6 6
R _I (34) ^c	37.6 102.6 112.3 498.0	T T	M W VW M	пе w 6 6 6
R _I (32)	131.6 ^a 165	1	M W	uem uem
R _I (30)	67.5 ^a 350	Ţ	M W	new 5
R _I (28)	310	T	W	5

- vs

Table I	(Continued)
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R _I (18)	41.5 ^a 43.7 ^a 219 858.3 ^a	# # #	VS VS W S	6 6 5,6
R _I (16)	81.6ª 86.4	1	W	6 6
R _I (14)	136.6 ^a 389.1	1	W	new 10
R _I (8)	41.5 71.0a 203.3 553 646.5a	1.d 1.d 1.d	M M VW - W	6 6 new 5 6
P _I (10)	108.7ª	H	W	10
P _I (18)	144.1 ^a 287.3 ^a 290	Ĭ	S S M	6 5,6 6
P _I (22)	34.8 40.1 258.4 ^a	# *	н н н	6 6 5,6
P _I (24)	238.3 286.7ª	# 	W M	6 5
P _I (28)	190 276.7ª	# 11	W M	new 5
P _I (32)	76.1 147.3 ^a 215.1 ^a	N N	M W M	6 10 new
P _I (42)	76.3 188.4 ^a	-	s s	new

^aFrequency measured

bSingle, double, and triple primes indicate different

pump offset frequencies. Cother $R_{\rm I}(34)$ pumped lines not observed in present experiment. Pump offset frequency not identified. Cour measured polarization is opposite to that of Ref. (6).

TABLE II. Frequency measurements for 39 FIR lasing transitions of CD_3OH (methanol-d₃) pumped by a CO_2 laser.

CD ₃ OH			ured	Vac				3OH	CO)
Laser Lin λ (μm)		requ HM)		Waven (cm	umt -1)	er ^D		r Press. ^C Pa)	Pump	Line
					·					
41.4	7	249	266.0	241.8	09	485	;	16	RT	(18)
43.7			664.2	228.8	47	124		16	RT	(18)
67.5	4	442	724.8	148.1	93	349	•	33	RT	(30)
71.0	4	223	062.0	140.8	66	186		17	RT	(8)
81.6	3	675	859.9	122.6	13	488		21	RT	(16)
86.7			161.2	115.2	85	129		12	RT	(34)
108.7	2	758	781.7	92.0	23	051	:	20	PT	(10)
120.7	2	484	584.9	82.8	76	831		40	RTI	(14)
122.2	2	454	225.9	81.8	64	164		47	RT	(38)
128.0	2	341	508.9	78.1	04	331		3	Rt	(34)
131.6	2	278	703.0	76.0	09	351	:	24	R _T	(32)
136.6	2	194	236.9	73.1	91	864		21	RT	(14)
144.1	2	080	189.3	69.3	87	647		41	PT	(18)
147.3	2	034	573.6	67.8			:	20	PT	(32)
159.4	1	880	754.6	62.7				16	RTT	(28)
168.1	1	783	601.1	59.4				4	R _T	(34)
180.7	1	658	689.9	55.3	27	940	:	21	RT	(34)
181.7	1	649	830.3	55.0				16	RTT	(28)
182.6	1	642	101.9	54.7	74	623		33	Rti	(14)
188.4	1	591	053.2	53.0				23	PT	(42)
189.7	1	580	101.8	52.7				40	Pti	(36)
191.4	1	566	672.8	52.2				23	R _T	(34)
198.7	1	508	908.6	50.3		774		36	Pty	(40)
215.1	1	39 3	856.9	46.4				21	P-	(32)
222.2			100.1	45.0		136	•	_	R _T	(34)
253.7	1	181	588.9	39.4			:	27	R+	(36)
258.4	1		027.8	38.6				23	P -	(22)
264.8			320.1	37.7				17		(34)
			395.1	36.1						\

Table II (Continued)

CD ₃ OH Laser Lin λ (μm)		Frequ	sured uency ^a Hz)	Vacuum Wavenum (cm ⁻¹)	er ^b	CD ₃ OH Laser Press (Pa)	CO ₂
286.2	1	047	502.3	34.940	916	20	P _{II} (40)
286.7	1	045	578.0	34.876	729	27	$P_{T}^{11}(24)$
287.3	1	043	454.5	34.805	896	27	P_{T}^{\perp} (18)
352.5		850	468.0	28.368	559	55	$R_{IT}^1(14)$
370.5		809	193.2	26.991	780	16	$R_{II}^{II}(28)$
386.0		776	589.1	25.904	223	13	$R_{T}^{11}(34)'$
418.7		715	987.6	23 882	776	15	R_{T}^{1} (36)
430.9		695	691.5	23.205	772		R _T (34)'
646.5		463	732.4	15.468	446	14	R_{T}^{\perp} (8)
858.3		349	305.1	11.651	566	11	R_T^1 (18)

The estimated fractional error in each measured frequency is ± 2x10⁻⁵.

b Calculated from the measured frequency with c = 299 792 458 m/s.

c 133.3 Pa = 1 torr.